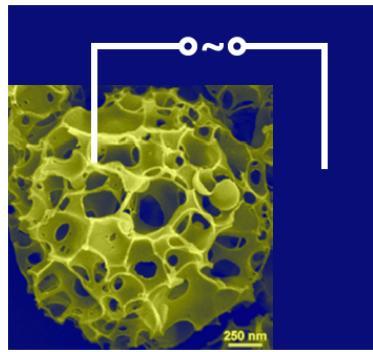


Ions at aqueous interfaces

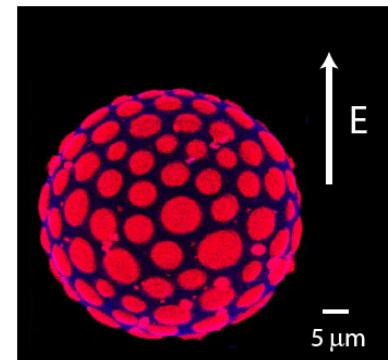
Douwe Jan Bonthuis

Ions adsorb at or desorb from surfaces

Double layer capacitance



Electrophoretic mobility
& conductivity

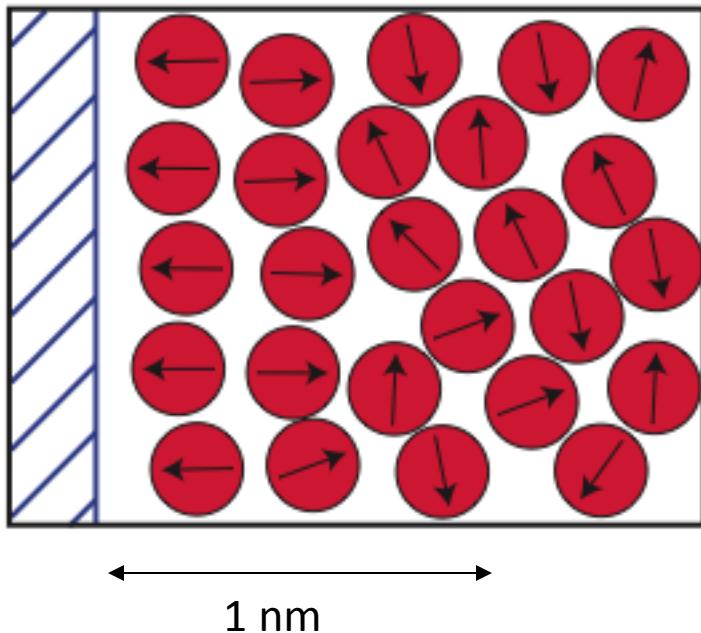


Colloidal stability
Gel electrophoresis
Vesicular transport in cells

Gouy-Chapman, DLVO do not
reproduce experimental results

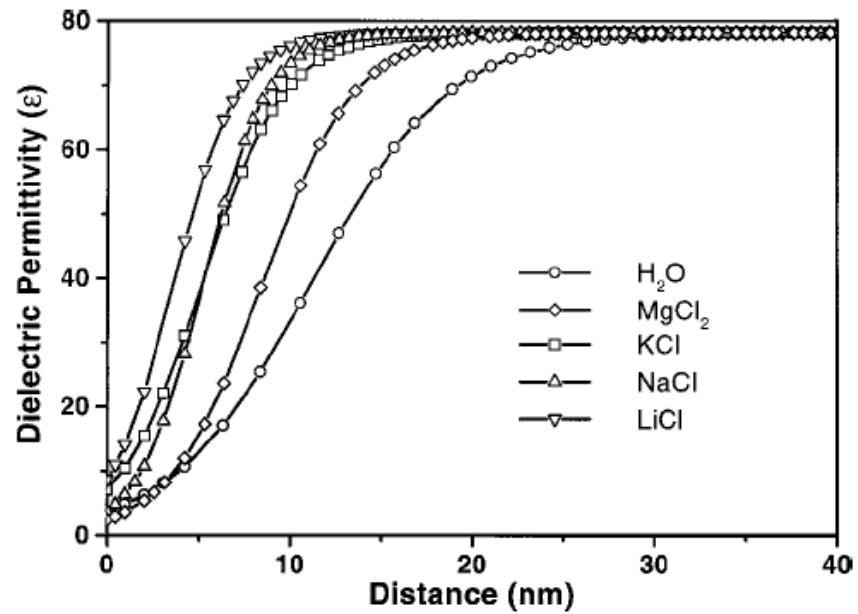
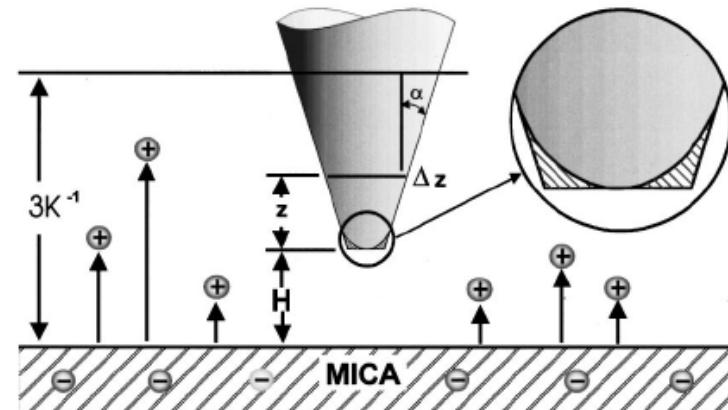
Interfacial Water Layers

Bulk



Orientation

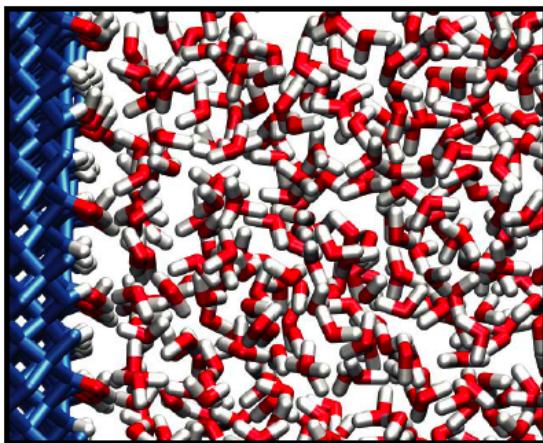
Structure



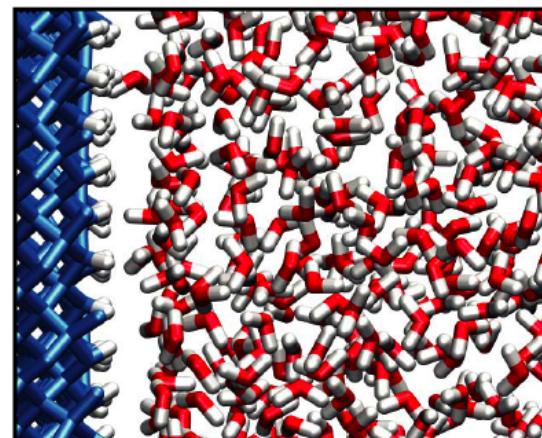
Teschke et al., PRE 64, 011605 (2001)

Molecular dynamics simulations at planar solid interfaces

Hydrophilic

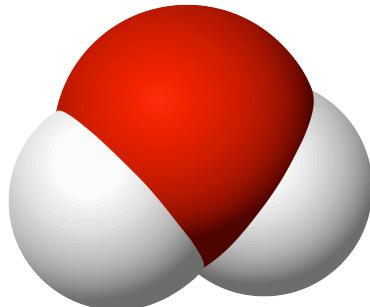


Hydrophobic



z →

SPC/E:



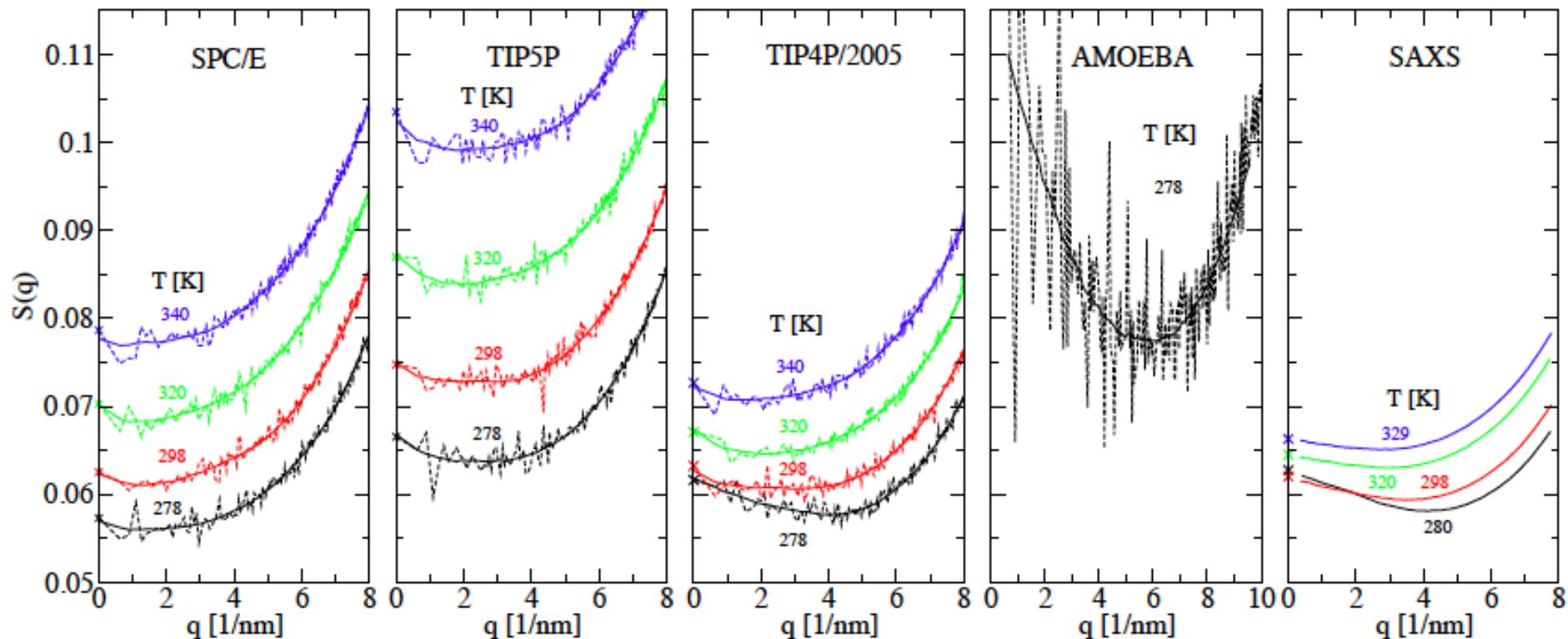
Rigid, non-polarizable

Permanent partial charge distribution

Spherically symmetric Lennard-Jones potential

SPC/E water (bulk)

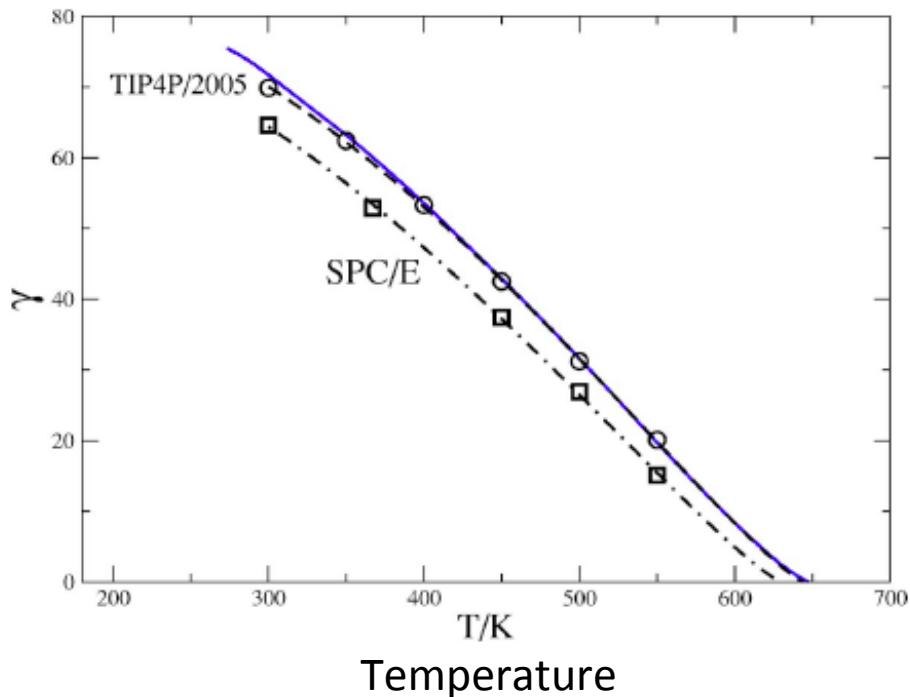
Structure factor



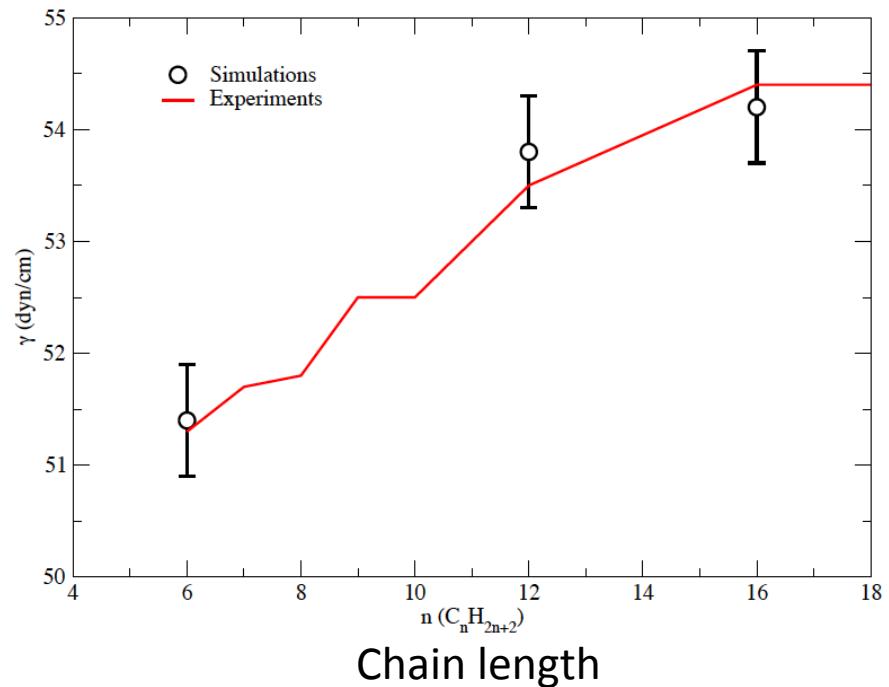
SPC/E water (interfaces)

Surface tension

Air/water



Oil/water



Vega & De Miguel,
J. Chem. Phys. **126**, 154707 (2007)

Dielectric tensor

$$\Delta D(r) = \int \varepsilon_0 \varepsilon_{nl}(r, r') \cdot \Delta E(r') dr'$$

$$\Delta E(r) = \int \varepsilon_0^{-1} \varepsilon_{nl}^{-1}(r, r') \cdot \Delta D(r') dr'$$

Local dielectric tensor at planar surfaces:

$$\varepsilon_{\parallel}(z) = \int \varepsilon_{nl}(z) dz$$

$$\varepsilon_{\perp}^{-1}(z) = \int \varepsilon_{nl}^{-1}(z) dz$$

Dielectric tensor

Excess polarization for a constant field \mathbf{F}

$$\Delta m = \langle m \rangle_F - \langle m \rangle_0$$

$$= \frac{\int (m - \langle m \rangle_0) \exp[-\beta(U - \mathbf{M} \cdot \mathbf{F})] dX}{\int \exp[-\beta(U - \mathbf{M} \cdot \mathbf{F})] dX}$$

$$\Delta m \approx \frac{\int (m - \langle m \rangle_0)(1 + \beta \mathbf{M} \cdot \mathbf{F}) \exp[-\beta U] dX}{\int \exp[-\beta U] dX}$$

Parallel (constant E)

$$\varepsilon_{\parallel}(z) \approx 1 + \varepsilon_0^{-1} \beta [\langle m_{\parallel}(z) M_{\parallel} \rangle_0 - \langle m_{\parallel}(z) \rangle_0 \langle M_{\parallel} \rangle_0]$$

Perpendicular (constant D)

$$\varepsilon_{\perp}^{-1}(z) \approx 1 - \varepsilon_0^{-1} \beta [\langle m_{\perp}(z) M_{\perp} \rangle_0 - \langle m_{\perp}(z) \rangle_0 \langle M_{\perp} \rangle_0]$$

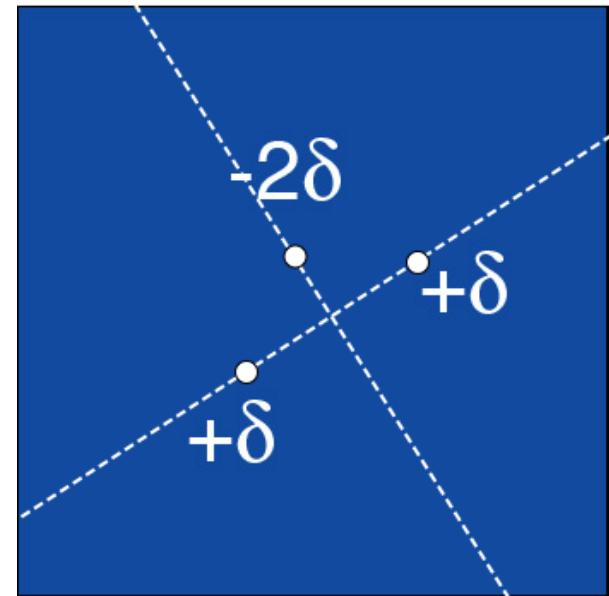
Expansion of the electric field

Directly from the charge density

$$m_\perp(z) = - \int_z^{\infty} \rho(z') \, dz'$$

$$m_{\parallel}(z) = \mp \int_V P_0(x, z) \, dx$$

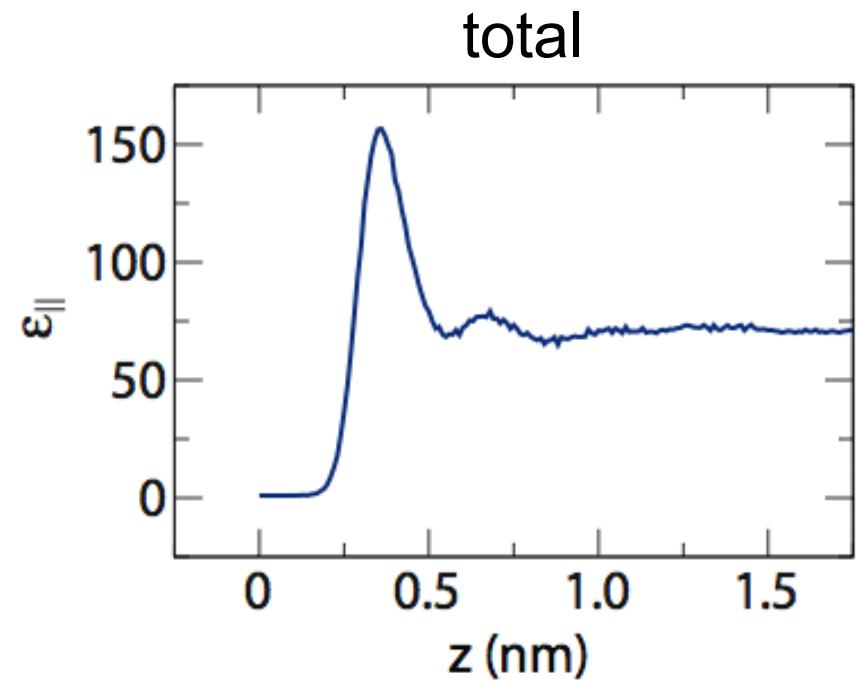
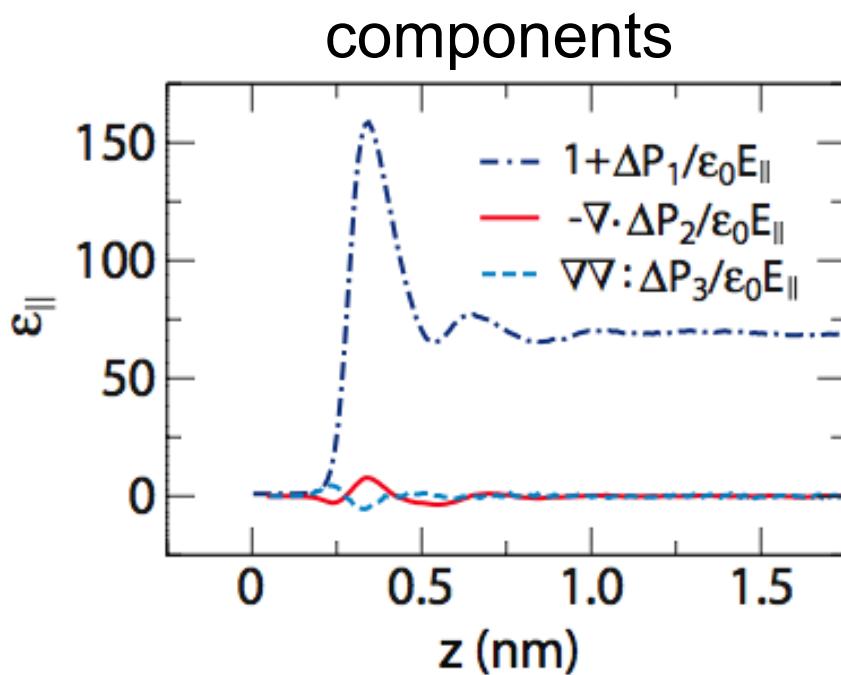
Expanding the polarization



$$m(r) = P_1(r) - \nabla \cdot P_2(r) + \nabla \nabla : P_3(r) - \dots$$

dipole quadrupole octupole

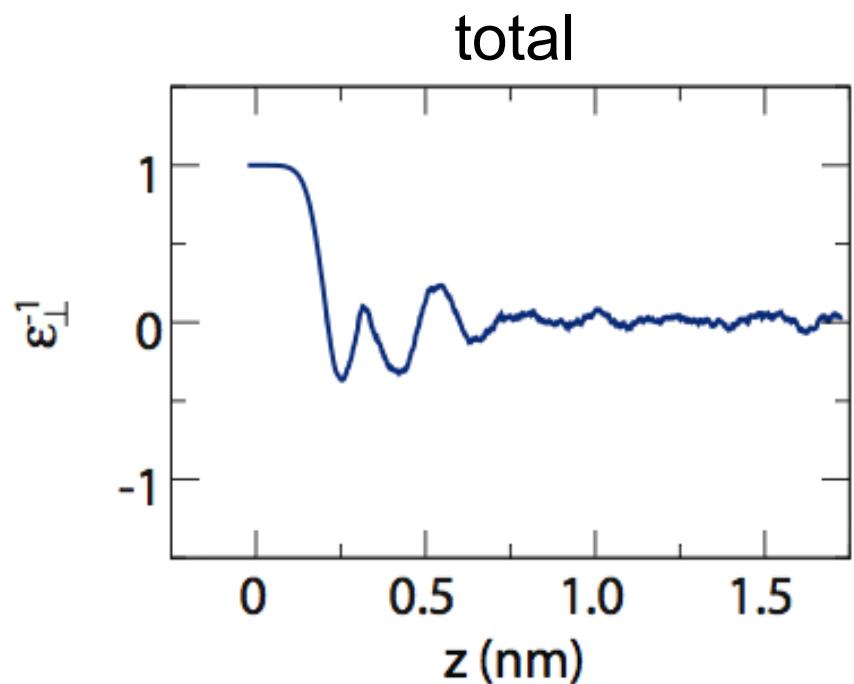
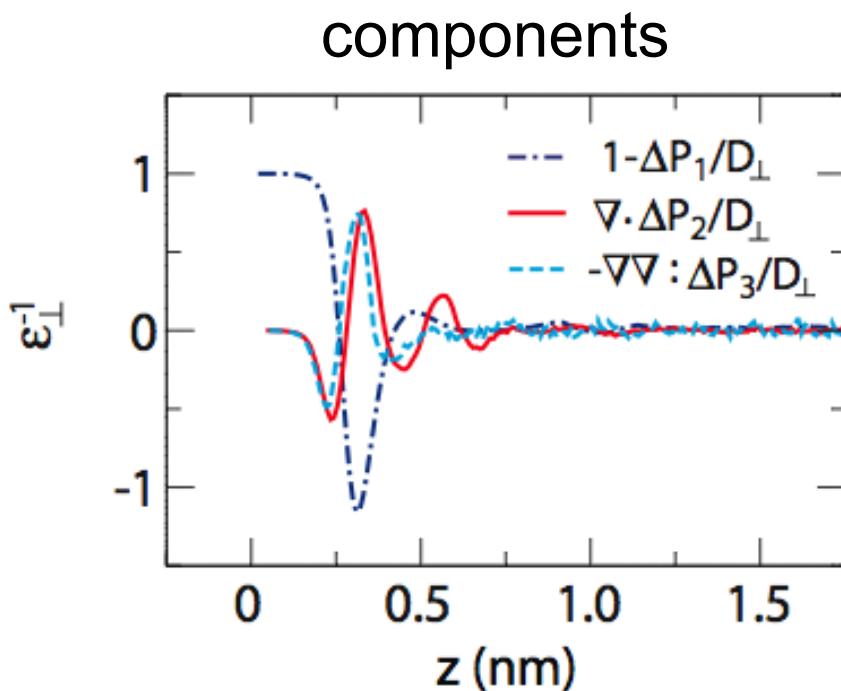
Parallel response



The dipole dominates the parallel response

Roughly proportional to the density

Perpendicular response



Significant contribution from higher order moments

$\epsilon_{\perp}^{-1} < 0$: overscreening

Extended Poisson-Boltzmann

$$\epsilon_0 E_\perp(z) = \epsilon_\perp^{-1}(z) D_\perp(z) \quad \text{Dielectric effects}$$

$$\epsilon_0 \nabla^2 \psi(z) = -\epsilon_\perp^{-1}(z) P_0(z) - D_\perp(z) \nabla \epsilon_\perp^{-1}(z)$$

$$D_\perp(z) = \int_0^z P_0(z') dz' \quad \text{Surface charges}$$

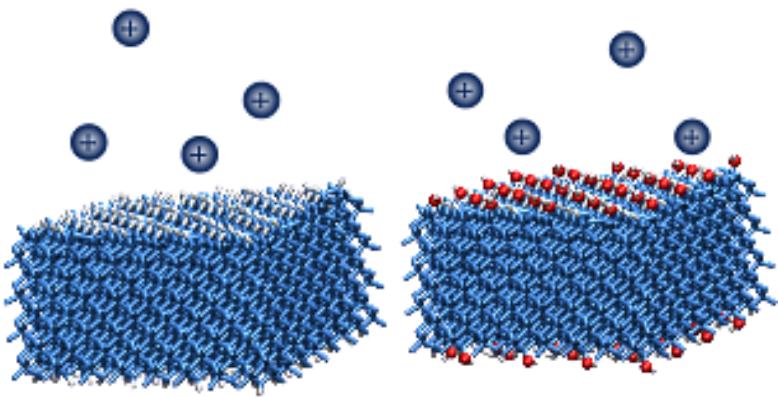
$$P_0(z) = e(c_+(z) - c_-(z)) \quad \text{Steric effects}$$

$$c_\pm(z) = \frac{\sqrt{2} \tilde{c}_\pm(z)}{\sqrt{2} + \lambda_+^3(\tilde{c}_+(z) - c_0) + \lambda_-^3(\tilde{c}_-(z) - c_0)}$$

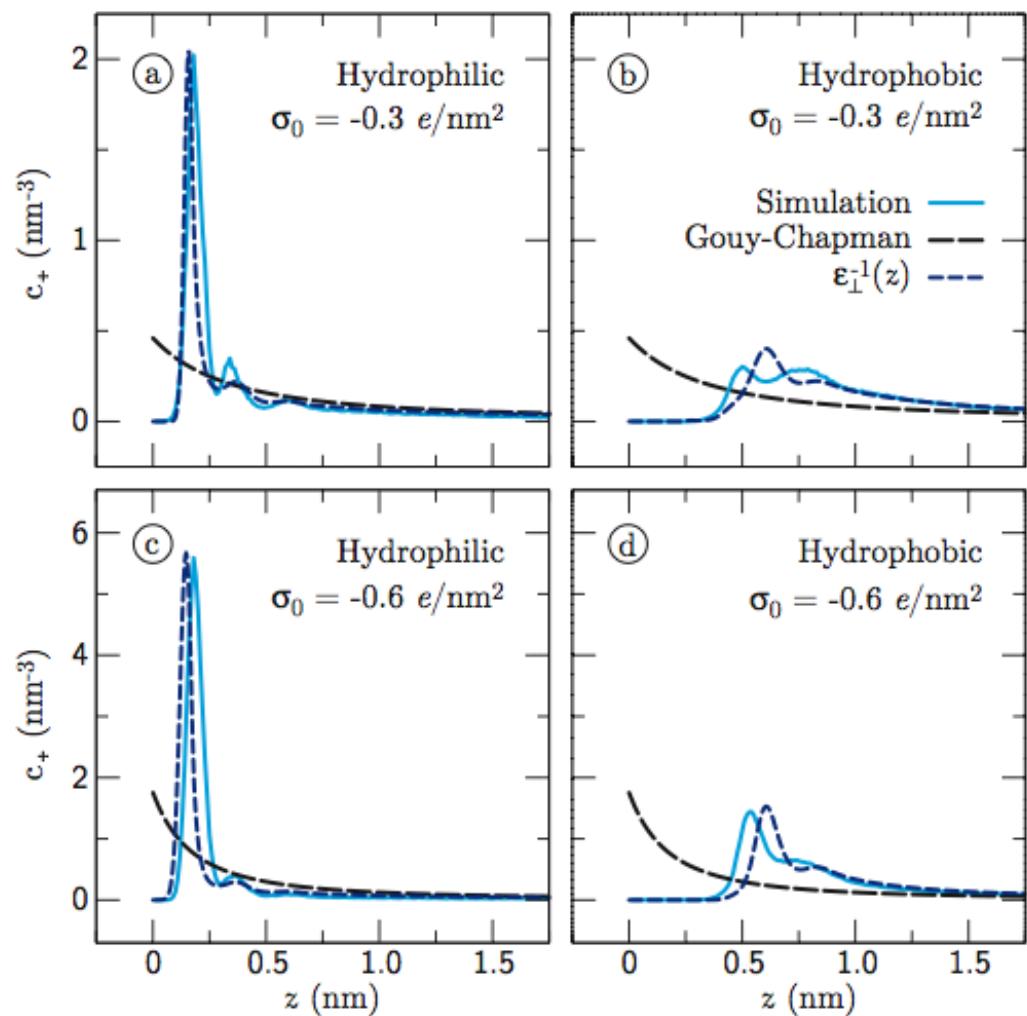
$$\tilde{c}_\pm(z) = c_0 \exp(-\mu_\pm(z) \mp \beta e \psi(z))$$

Other nonelectrostatic contributions

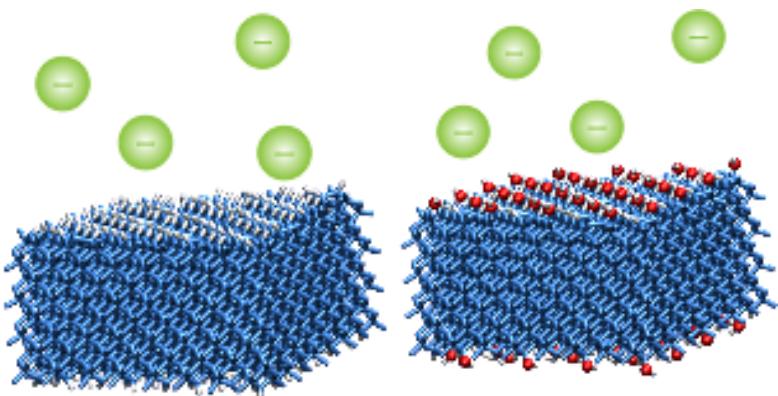
Sodium ion density



soft wall repulsion

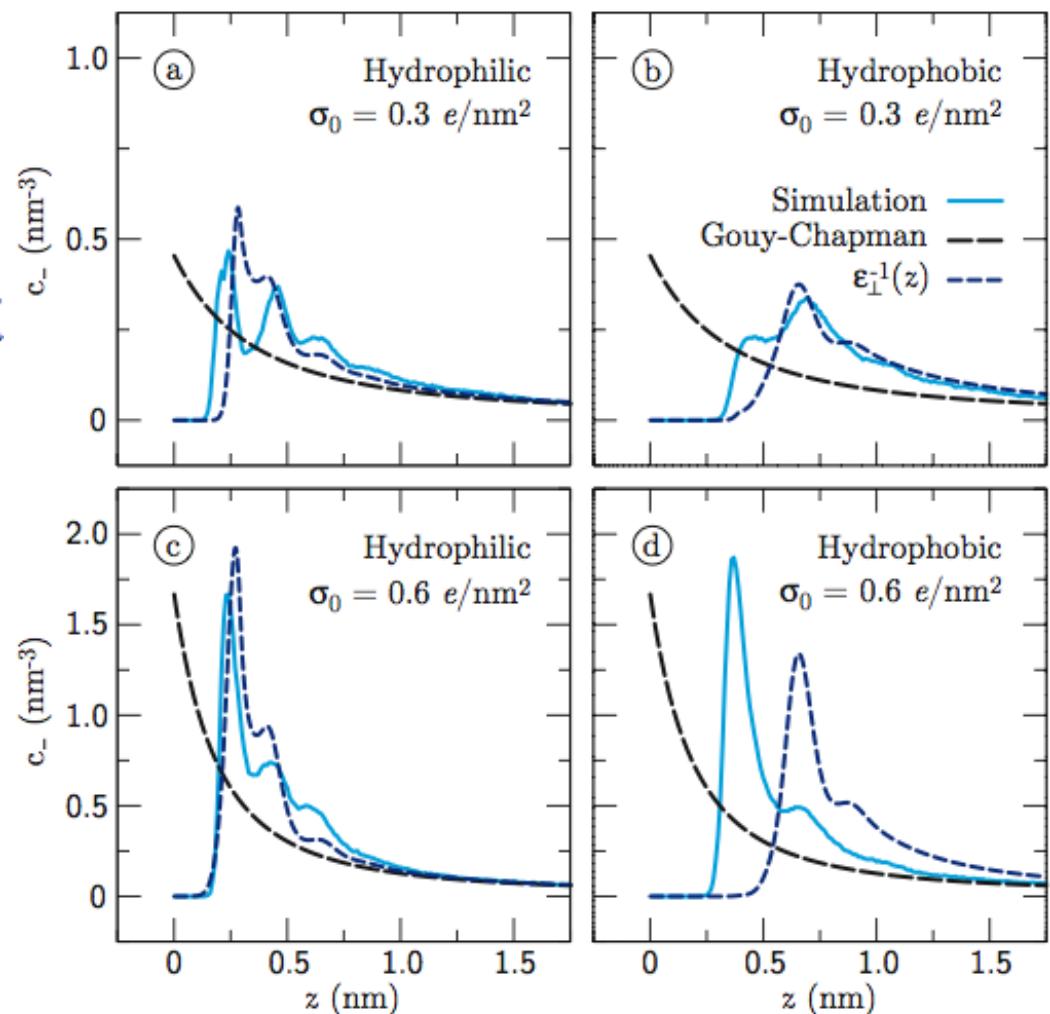


Chloride ion density

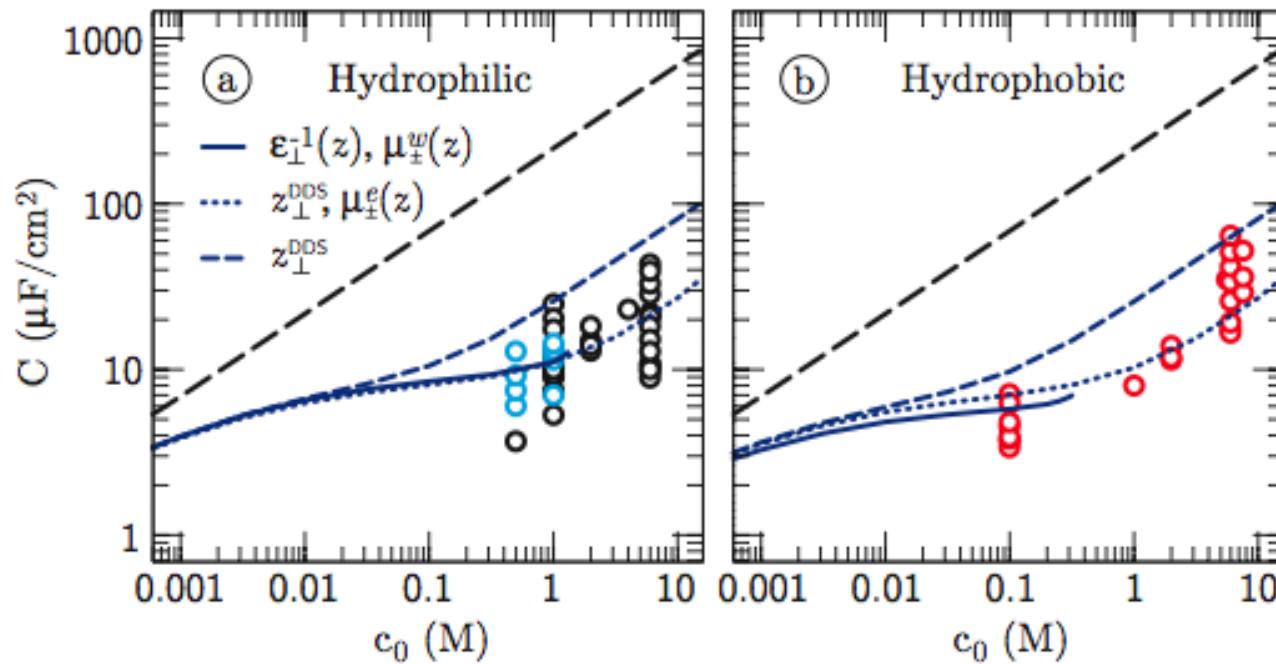


Dielectric profile & soft wall
reproduce ion density profiles

Except for big ions at
hydrophobic surfaces

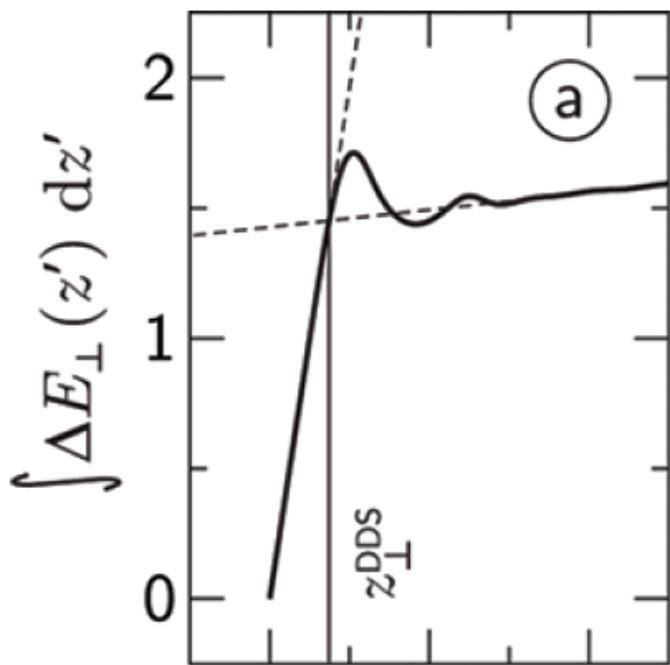


Double-layer capacitance



$$\mu_{\pm}^e(z) = \alpha \exp [1 - 2z/\lambda_{\pm}]$$

Analytical approximation



$$\varepsilon_{\perp}^{-1}(z) = \begin{cases} 1 & \text{if } z < z_{\perp}^{\text{DDS}} \\ \varepsilon_{\text{bulk}}^{-1} & \text{otherwise,} \end{cases}$$

Debye-Hückel:

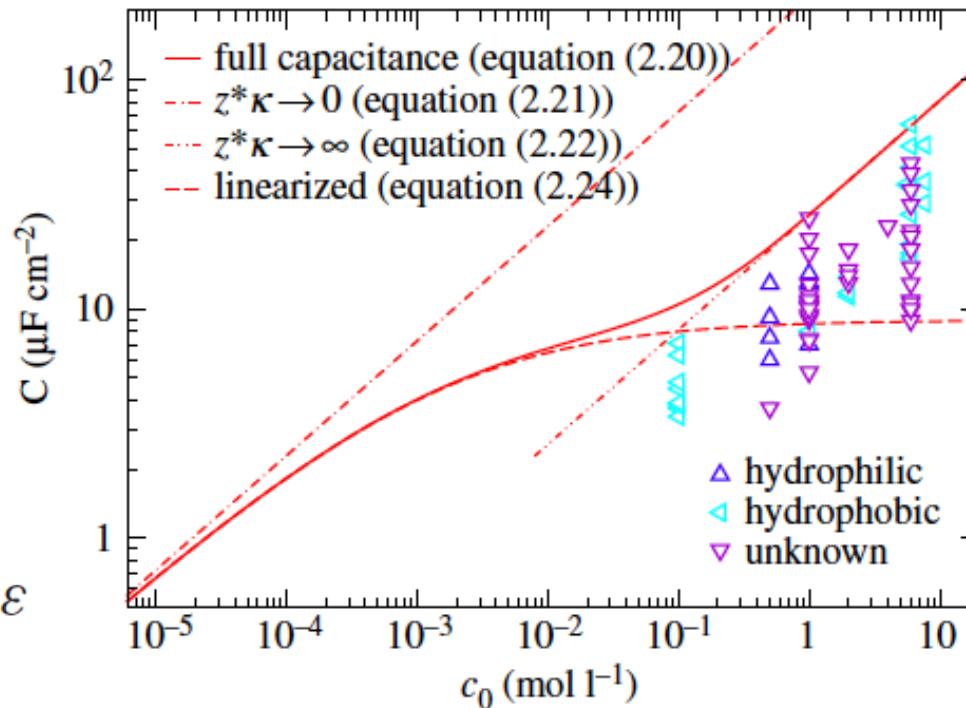
$$\nabla^2 \psi(z) = \begin{cases} \varepsilon \kappa^2 \psi(z), & \text{for } z < z^* \\ \kappa^2 \psi(z), & \text{for } z > z^* \end{cases}$$

$$C = \left(\frac{d\psi(z=0)}{d\sigma_0} \right)^{-1}$$

Capacitance

$$C = \varepsilon_0 \kappa \sqrt{\varepsilon} \frac{1 + e^{-2z^* \kappa \sqrt{\varepsilon}} (\sqrt{\varepsilon} - 1) / (\sqrt{\varepsilon} + 1)}{1 - e^{-2z^* \kappa \sqrt{\varepsilon}} (\sqrt{\varepsilon} - 1) / (\sqrt{\varepsilon} + 1)}$$

$$C \simeq \varepsilon_0 \kappa \varepsilon$$



$$C \simeq \varepsilon_0 \kappa \sqrt{\varepsilon}$$

To linear order in $z^* \kappa$

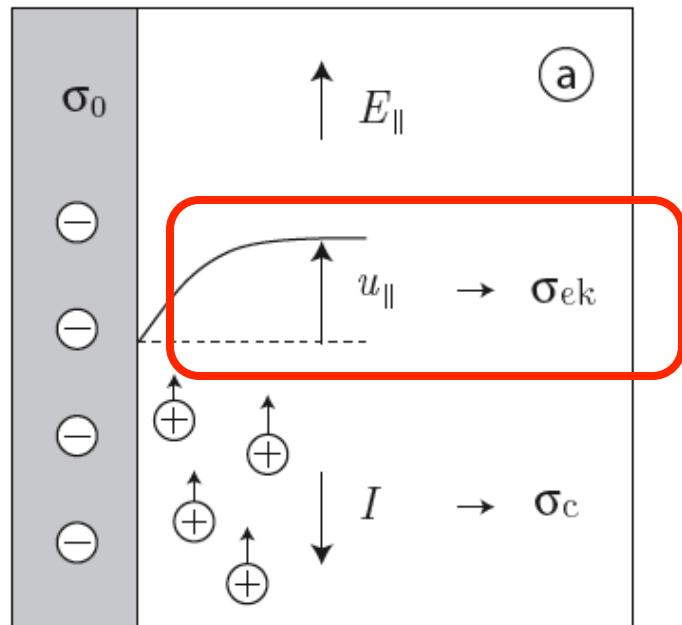
$$C \simeq \left(\frac{\kappa^{-1} - z^*}{\varepsilon_0 \varepsilon} + \frac{z^*}{\varepsilon_0} \right)^{-1}$$

2 capacitors in series!

Electrokinetics

Electroosmotic mobility

$$\lim_{z \rightarrow \infty} \frac{u_{\parallel}(z)}{E_{\parallel}} = -\frac{\varepsilon_0 \varepsilon_{bulk}}{\eta_{bulk}} \zeta$$



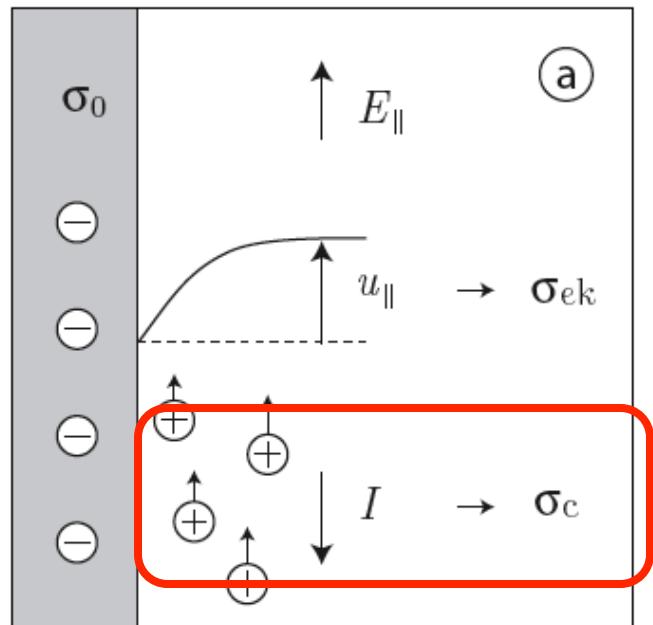
Electrokinetic surface charge density:

$$\sigma_{ek} = \sqrt{\frac{8c_0 \varepsilon_0 \varepsilon_{bulk}}{\beta}} \sinh \left[\frac{\beta e \zeta}{2} \right]$$

Electrokinetics

Surface conductivity

$$\frac{I}{E_{\parallel}} = \int_0^{\infty} e [c_+(z) - c_-(z)] [u_{\parallel}(z)/E_{\parallel}] dz + \int_0^{\infty} e [\nu_+ (c_+(z) - c_0) + \nu_- (c_-(z) - c_0)] dz$$

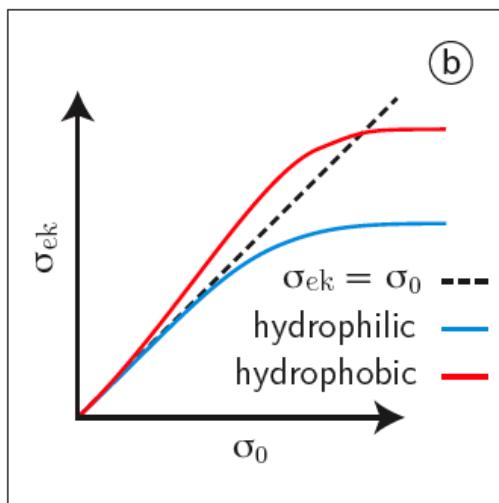


Conductive surface charge density:

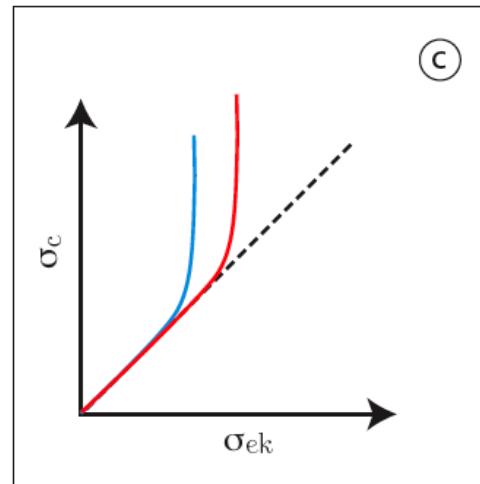
$$\sigma_c = \frac{\eta \kappa^2}{4ec_0 + \eta \nu \kappa^2} \sqrt{\frac{I}{E_{\parallel}}} \sqrt{\frac{I}{E_{\parallel}} + \frac{32e^2 c_0^2}{\eta \kappa^3} + \frac{8ec_0 \nu}{\kappa}}$$

Experimental inconsistencies

1. mobility



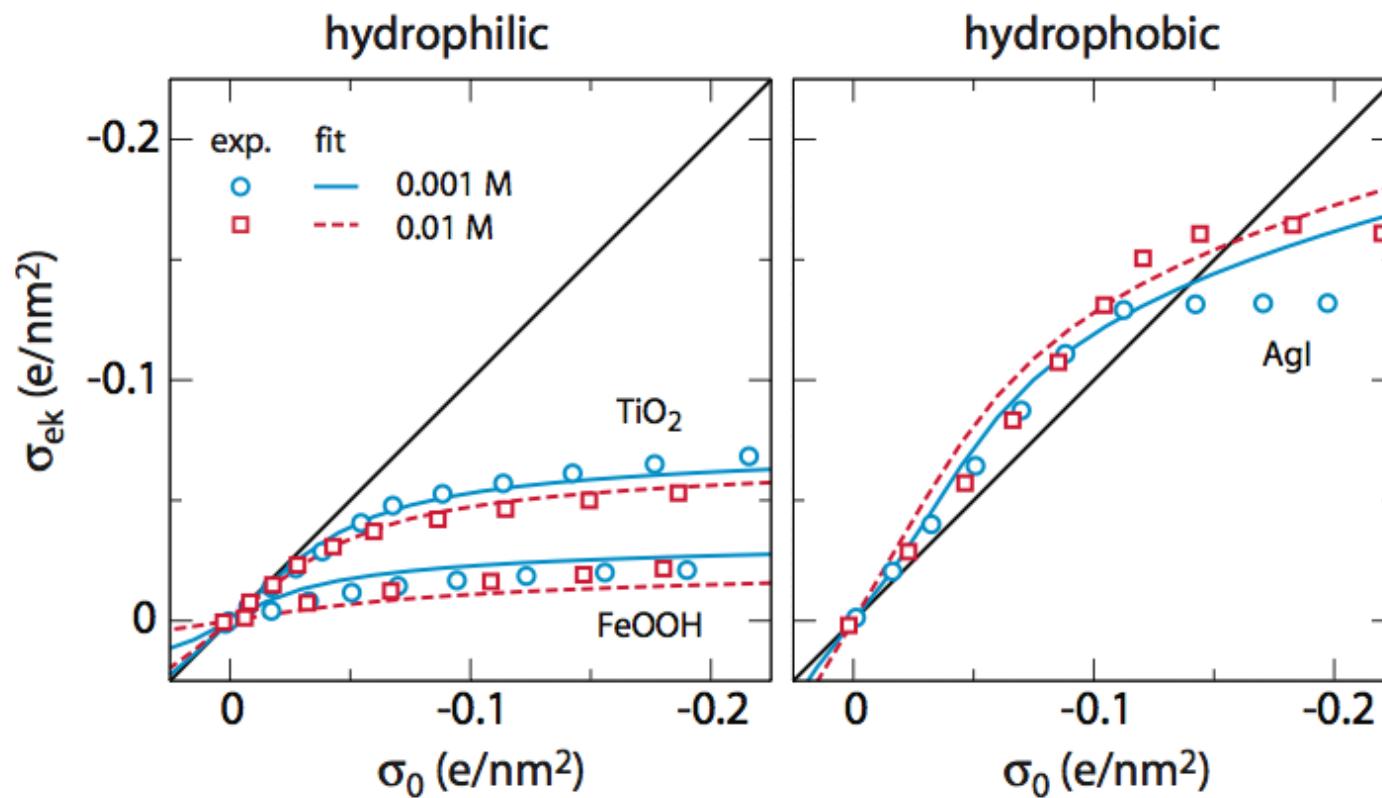
2. conductivity



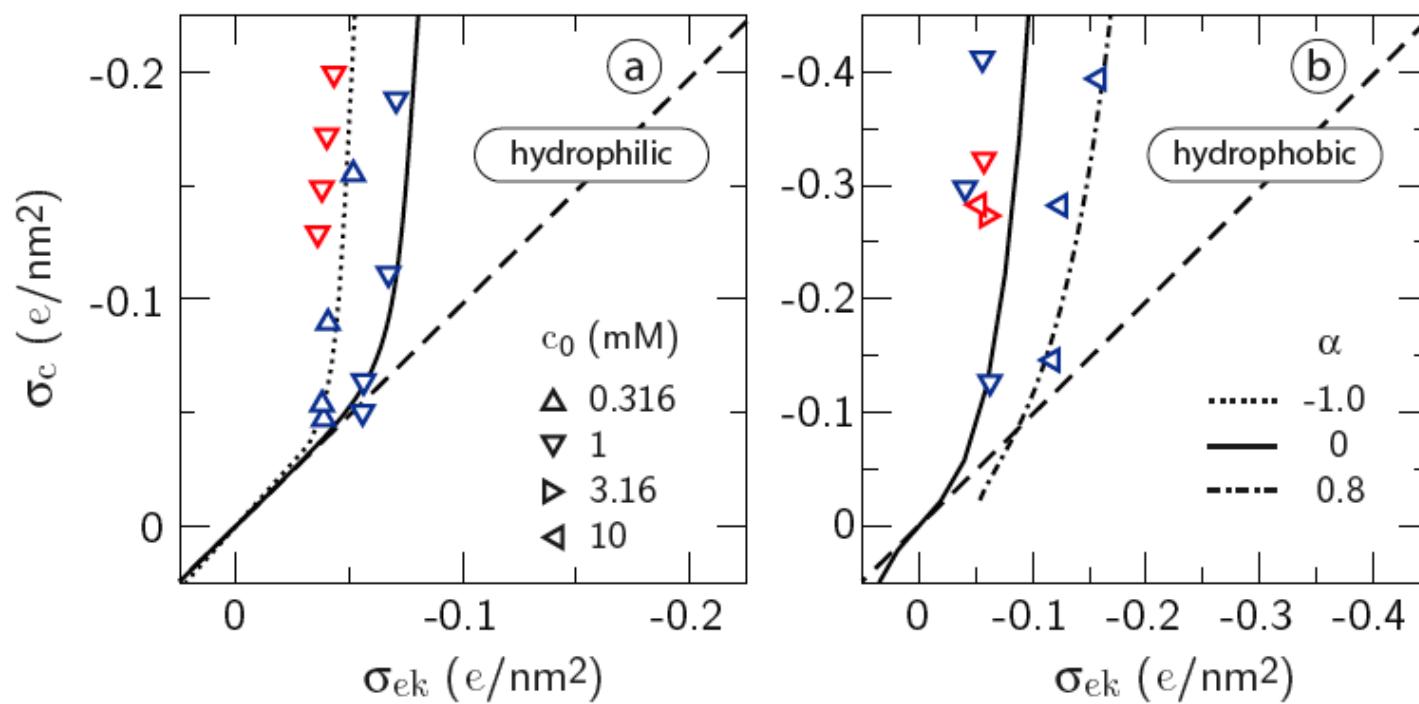
Electrophoretic mobility

Fit parameter:

$$\mu(z) = \alpha \exp[1 - z/\lambda]$$



Excess conductivity



Water ions

Auto-ionization

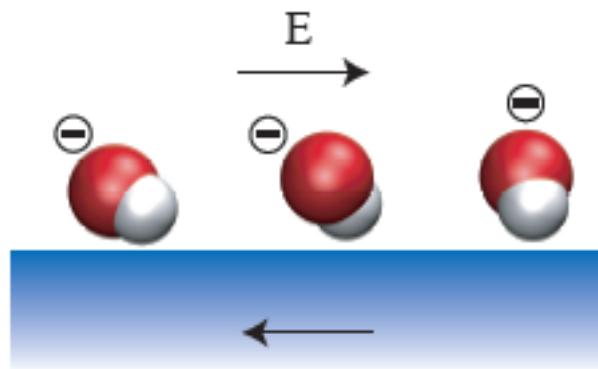


Always present, typically at very low concentration

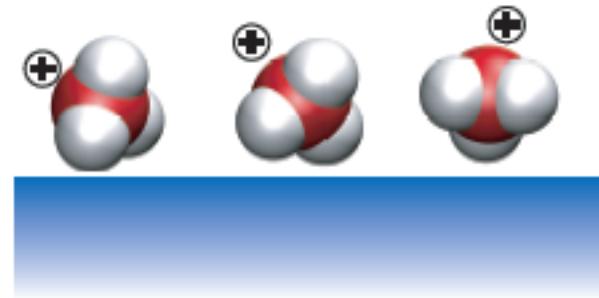
Experiments on bubbles and oil droplets

Electrophoresis
Beattie's sonication experiment

Spectroscopy
Second harmonics generation
Surface tension

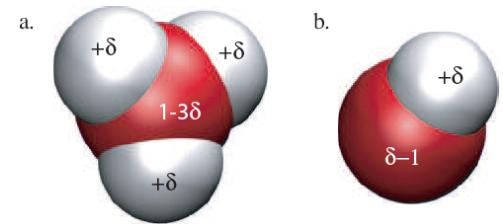
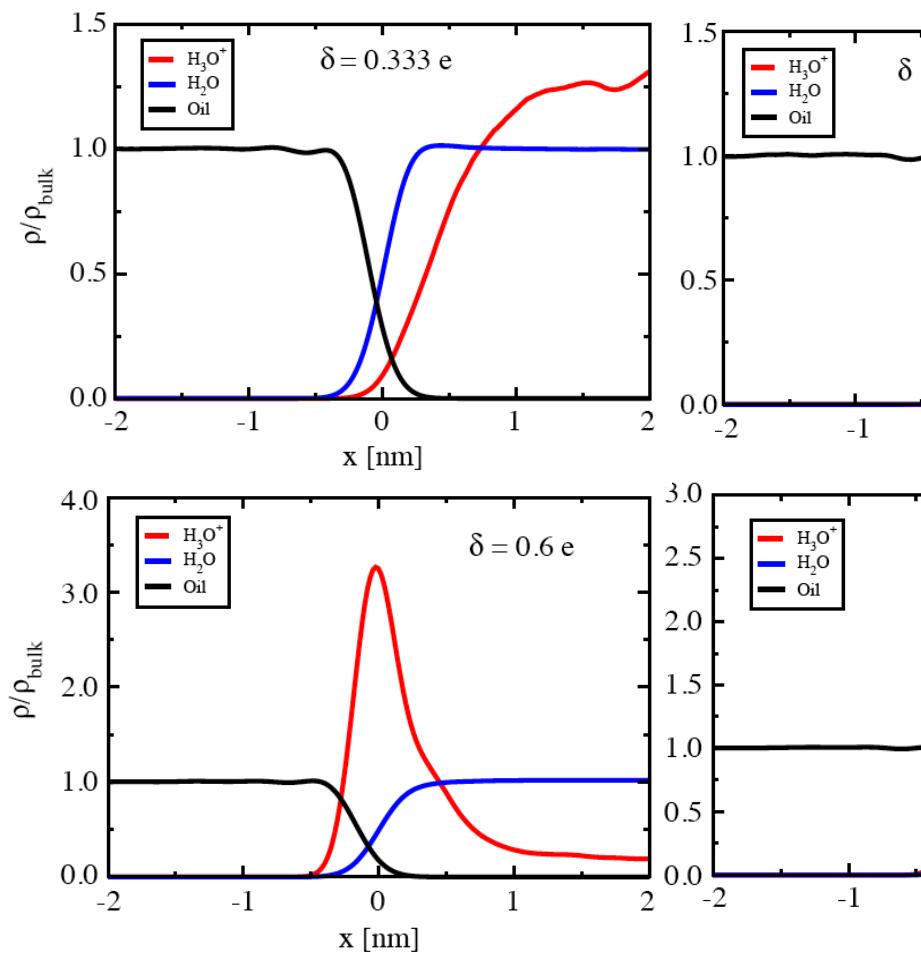


Negative effective surface charge

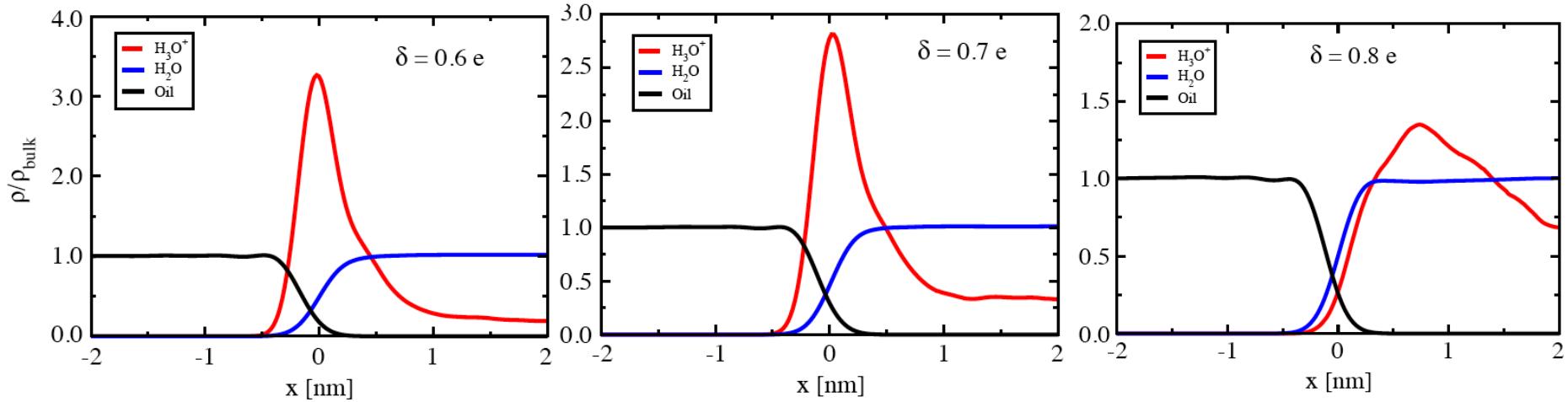


Positive effective surface charge

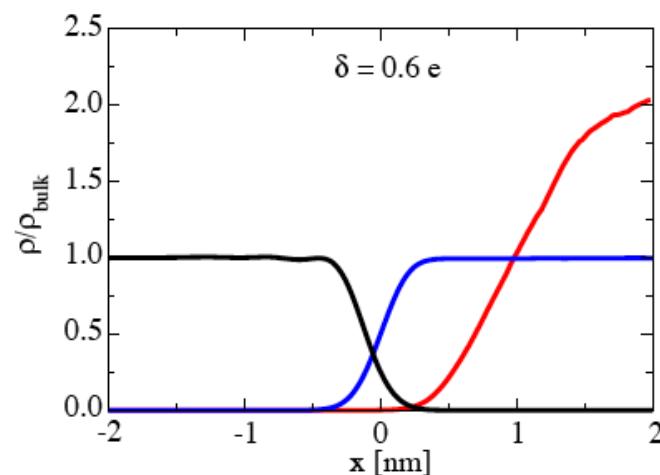
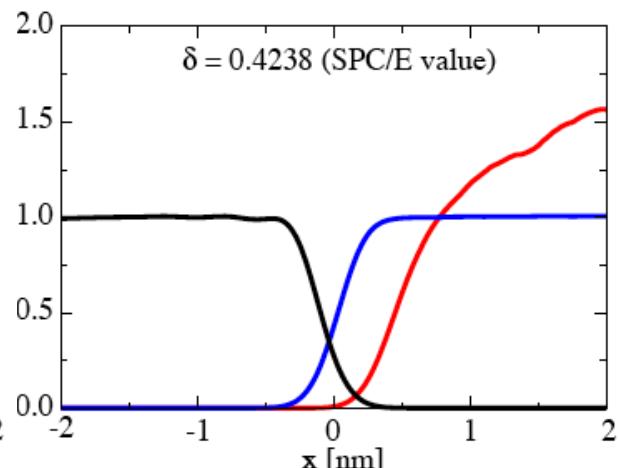
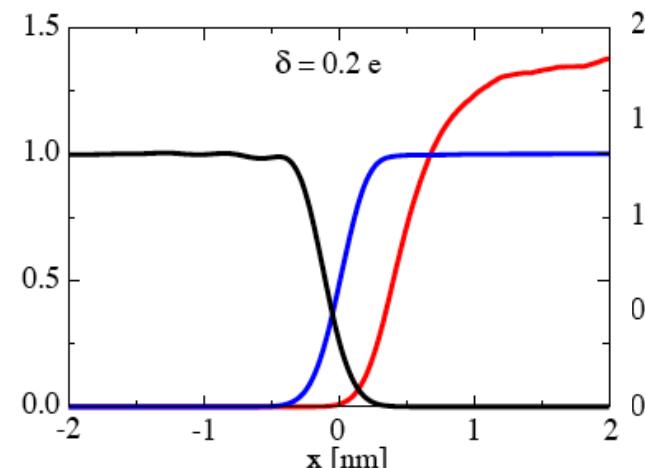
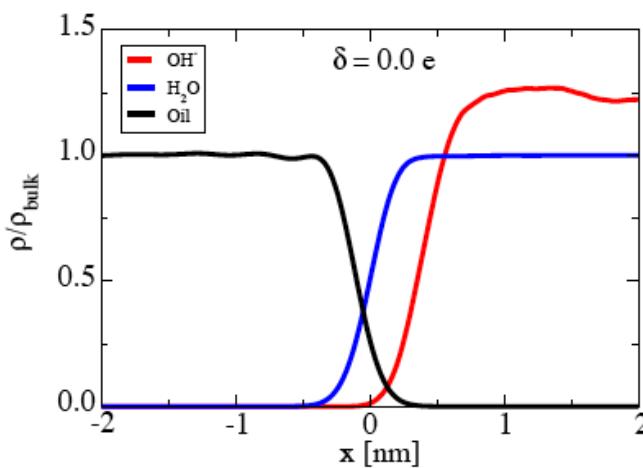
Effect of the charge distribution



The same Lennard-Jones parameters as SPC/E water



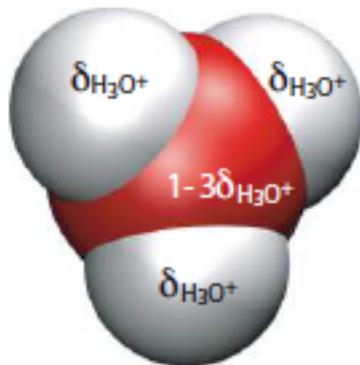
Effect of the charge distribution



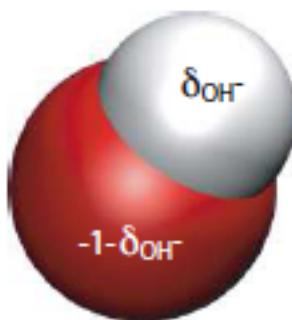
For both H_3O^+ and OH^- , the partial charges have a strong influence on the adsorption

Force fields

a)



b)



Like the water:
No atomic polarizability
No molecular polarizability

Parameters:

1. Lennard-Jones of the oxygen
2. Partial charge on the hydrogen
3. Combination rules

Observables:

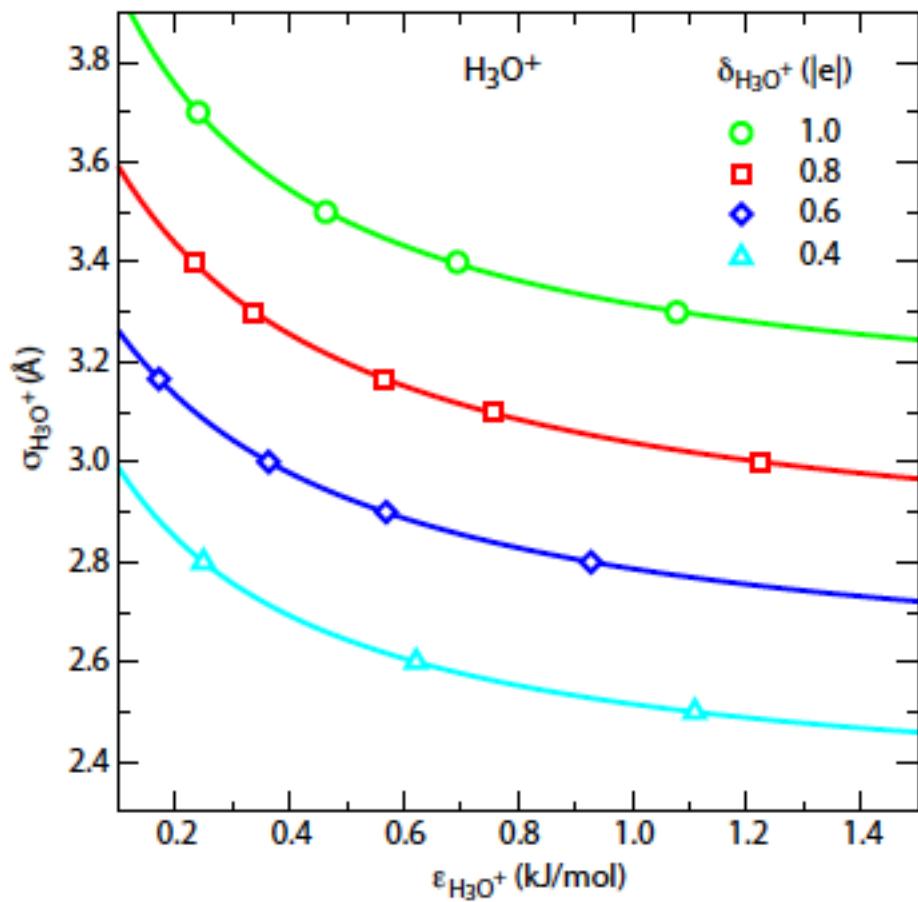
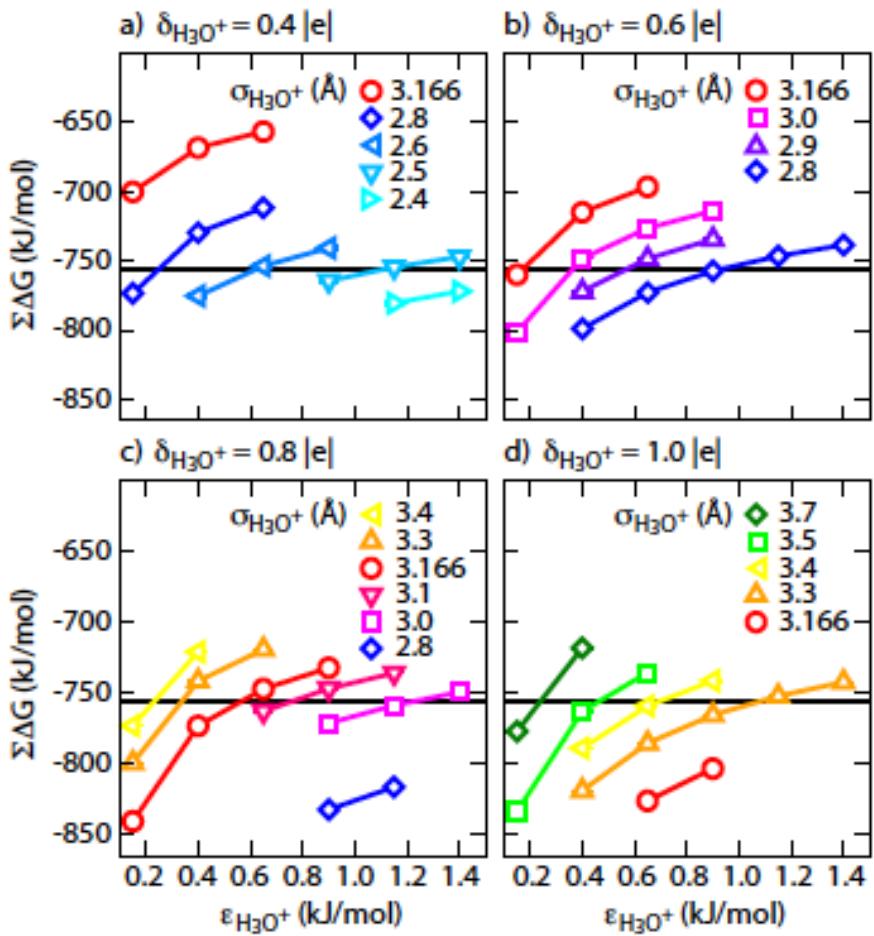
1. Ion-water interaction: solvation free energy
2. Ion-ion interaction: activity coefficient

$$\gamma_i(x_i)x_i = e^{-\beta(\mu - \mu_0)}$$

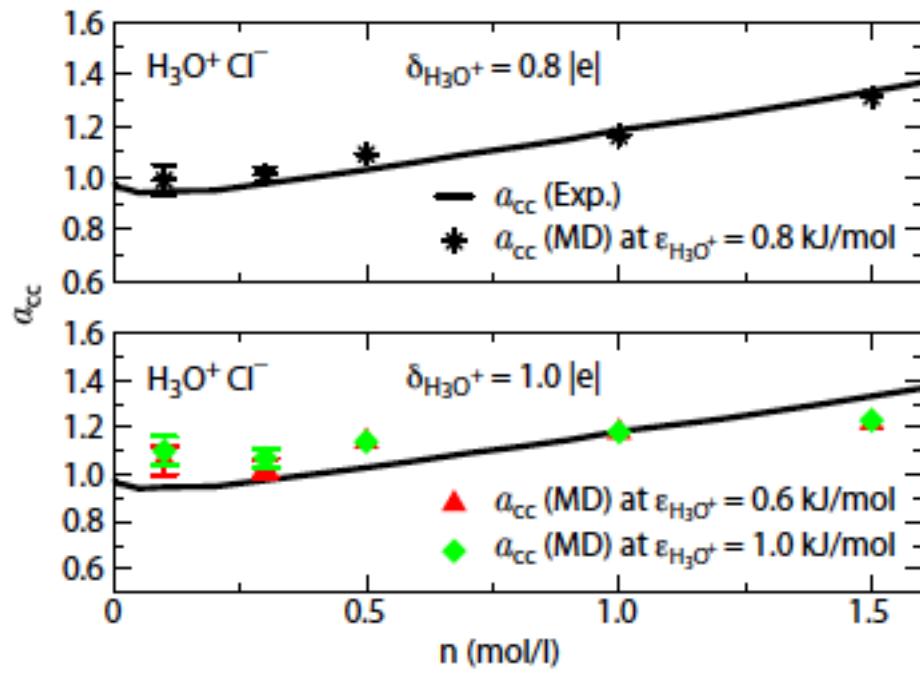
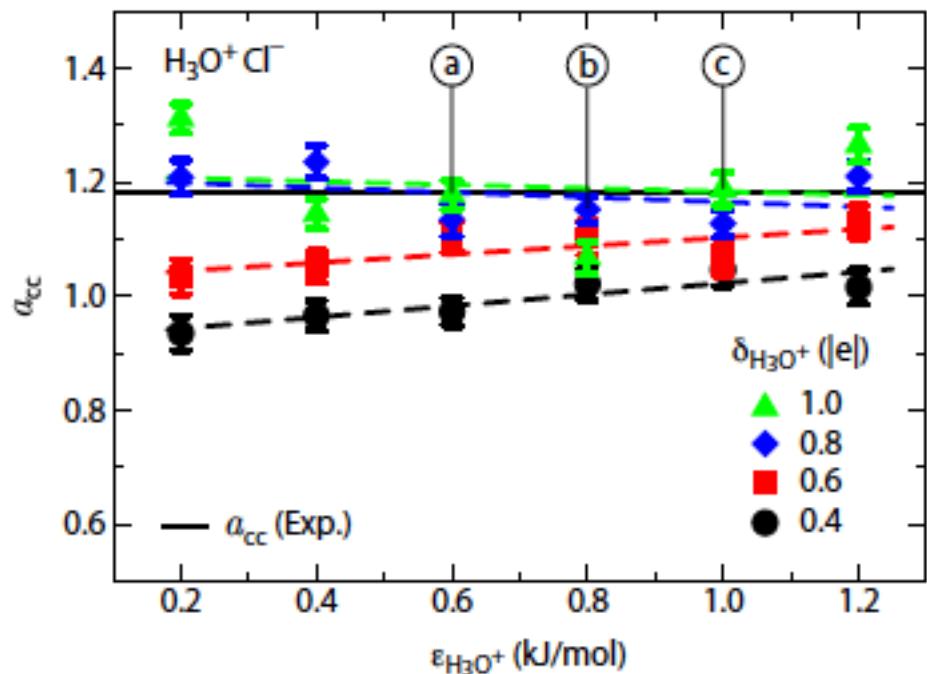
$$1 + \frac{\partial \ln \gamma}{\partial \ln n} = \frac{G_{+-}^\infty - G_{++}^\infty}{2(G_{+-}^\infty - G_{+s}^\infty)}$$

$$G_{\alpha\beta}^R(R) = \frac{1}{v(R)} \int_{v(R)} \int_{v(R)} [g_{\alpha\beta}(\mathbf{r}_1, \mathbf{r}_2) - 1] d\mathbf{r}_1 d\mathbf{r}_2$$

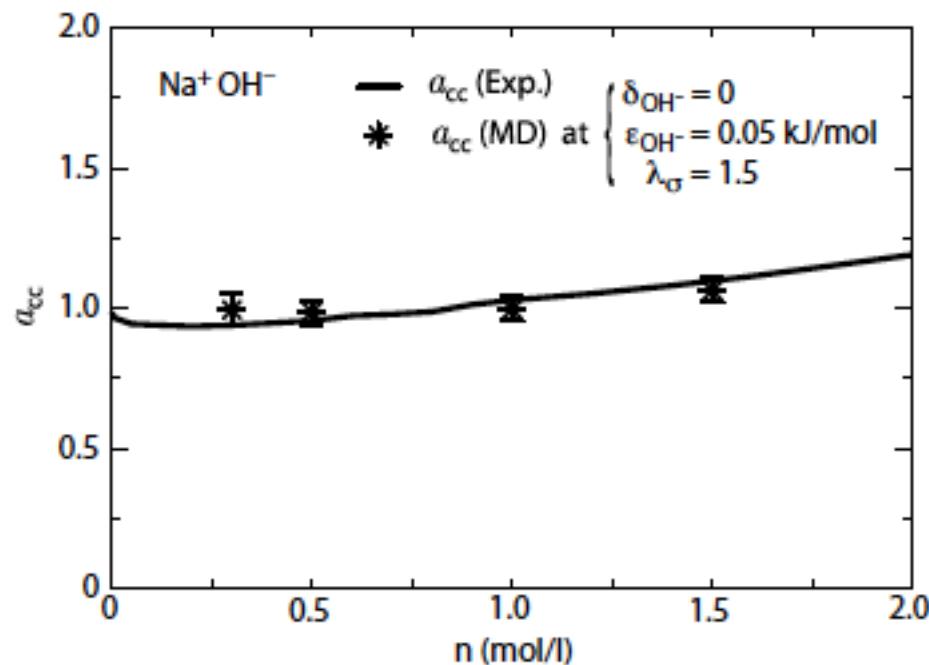
Free energy optimization



Activity optimization HCl

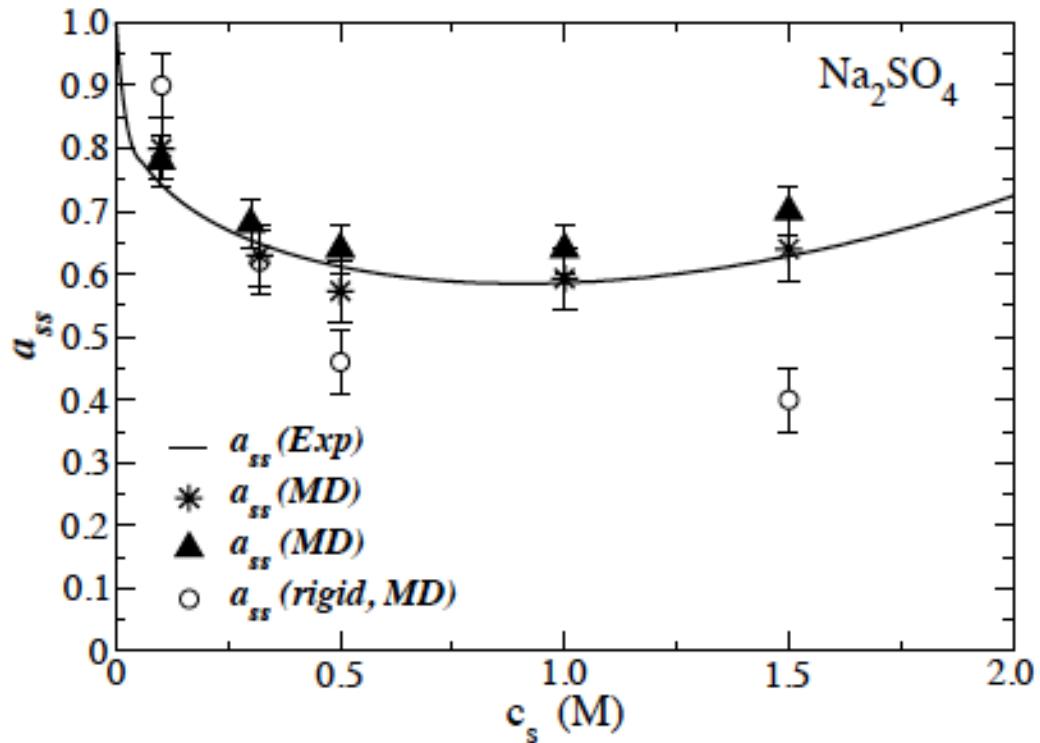
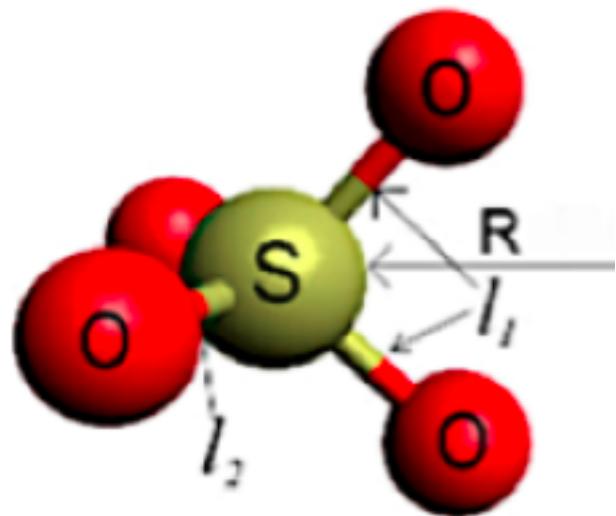


Activity optimization NaOH



$$\sigma_{+-} = \lambda_\sigma \sqrt{\sigma_+ \sigma_-}$$

Suphate



Slightly different procedure:

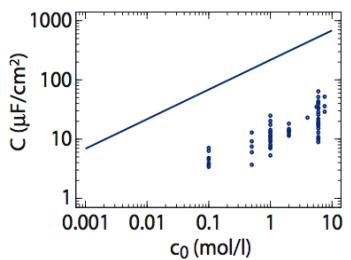
- Fix the partial charge
- Allow for some flexibility

Extra parameters:

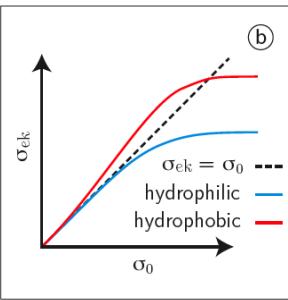
- Bonds
- Angles
- Intramolecular interactions

Conclusions

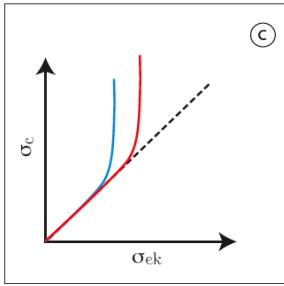
1. capacitance



2. mobility



3. conductivity



4. Optimization of thermodynamically consistent ion force fields

Thanks

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Shavkat Mamatkulov
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Lyderic Bocquet
Alexander Schlaich
Stephan Gekle

Glasstone Fellowship



Linacre
College

DFG

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